

Magnetic Field Dependence of the Sommerfeld Coefficient in $\text{Mg}_x\text{Al}_{1-x}\text{B}_2$ Single Crystals

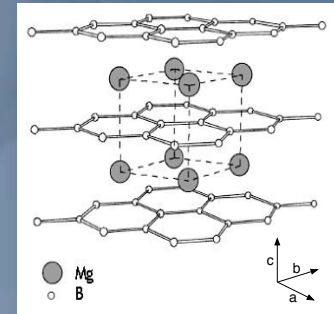
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Z.Holanova
C.Marcenat

Collaboration :
P. Samuely & P. Szabo : Point Contact Spectroscopy
S-I.Lee and coll. & S.Lee : sample preparation

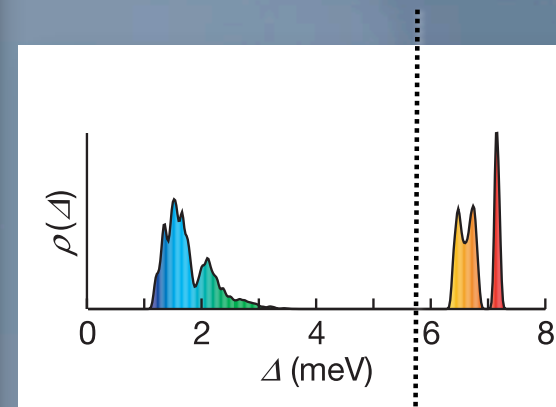
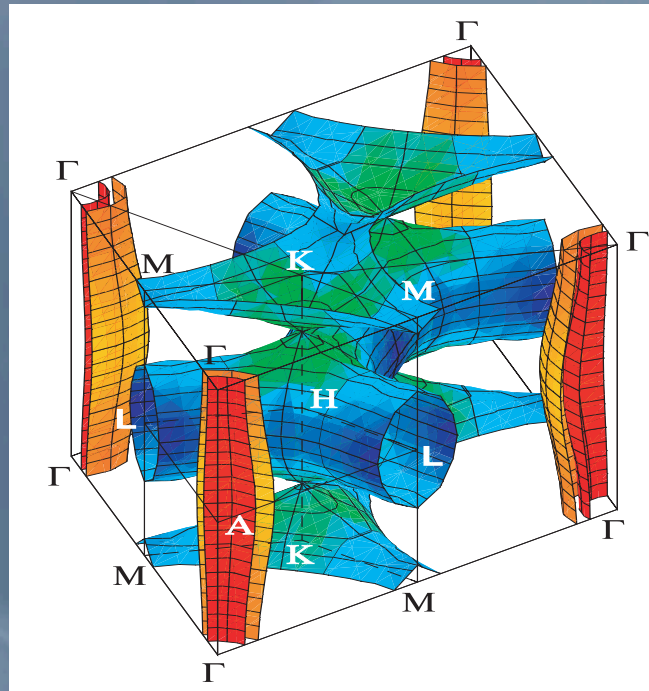
Hexagonal structure (B planes)

- quasi 2D sp^2 B orbitals : σ band
- 3D p_z B orbitals : π band

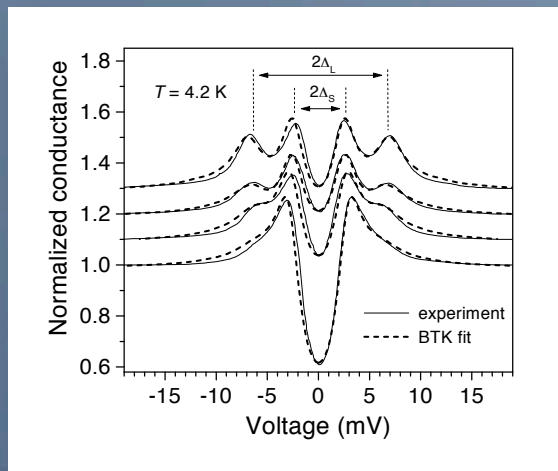
2 gap superconductor



$$T_c \sim 40K$$

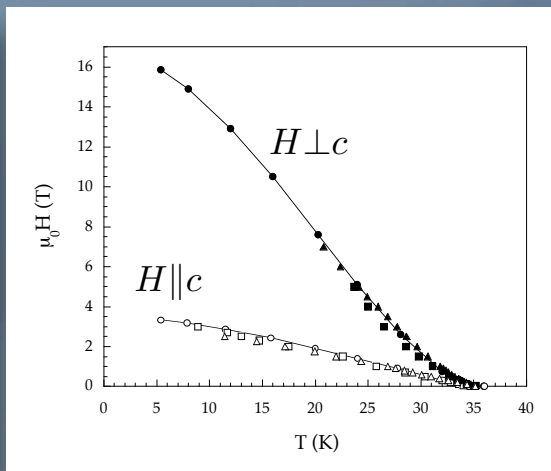


BCS value



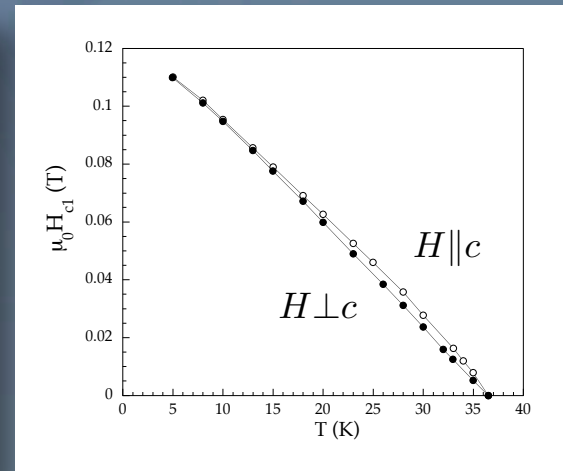
Tunneling spectroscopy
(and specific heat : bulk)

Rapid suppression of the small
gap with field



C_p : thermodynamic criterion for
the determination of $H_{c2}(T)$
(also $R=0$, onset of diamagnetism)

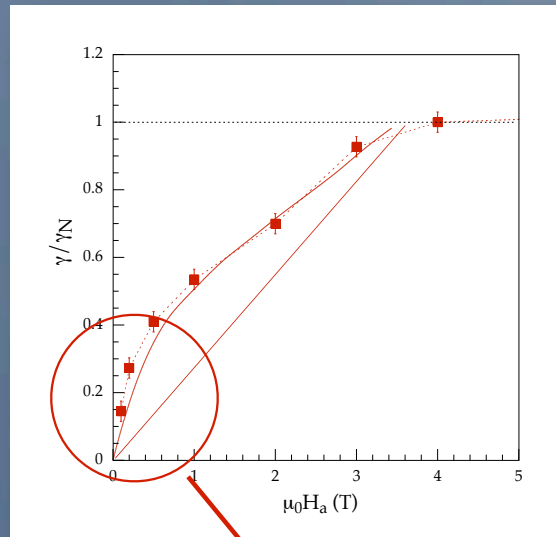
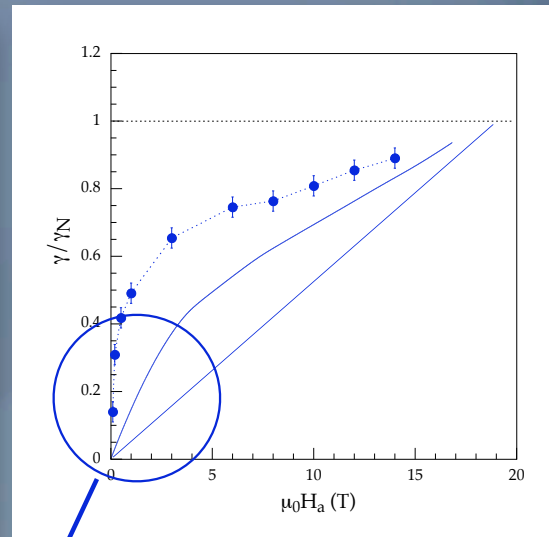
Temperature dependence
of the anisotropy



Magnetic measurements :
determination of $H_{c1}(T)$

isotropic $H_{c1}(T)$

$$\Gamma_{H_{c2}} \neq \Gamma_{H_{c1}}$$

$H \parallel c$

 $H \perp c$


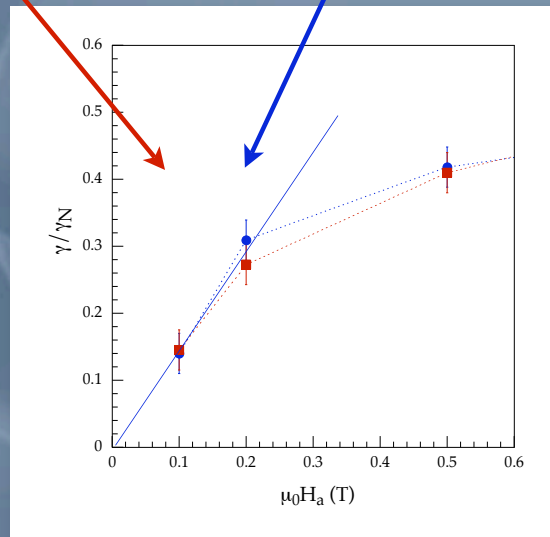
$$\gamma = \frac{C_p}{T} \Big|_{T \rightarrow 0}$$

Classical (dirty) superconductors

$$\gamma \propto \left\{ \frac{\xi}{a_0} \right\}^2 \propto \frac{B}{B_{c2}}$$

non-linear dependence predicted
by Kogan et al. due to core
shrinking in clean systems

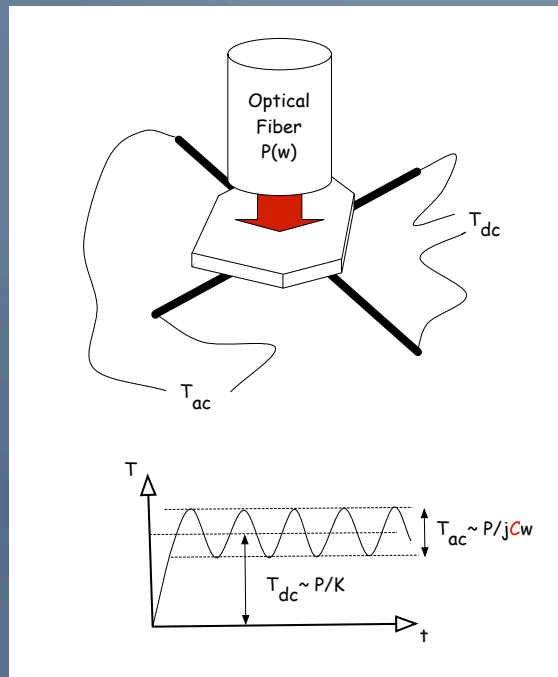
F.Bouquet et al. PRL 02



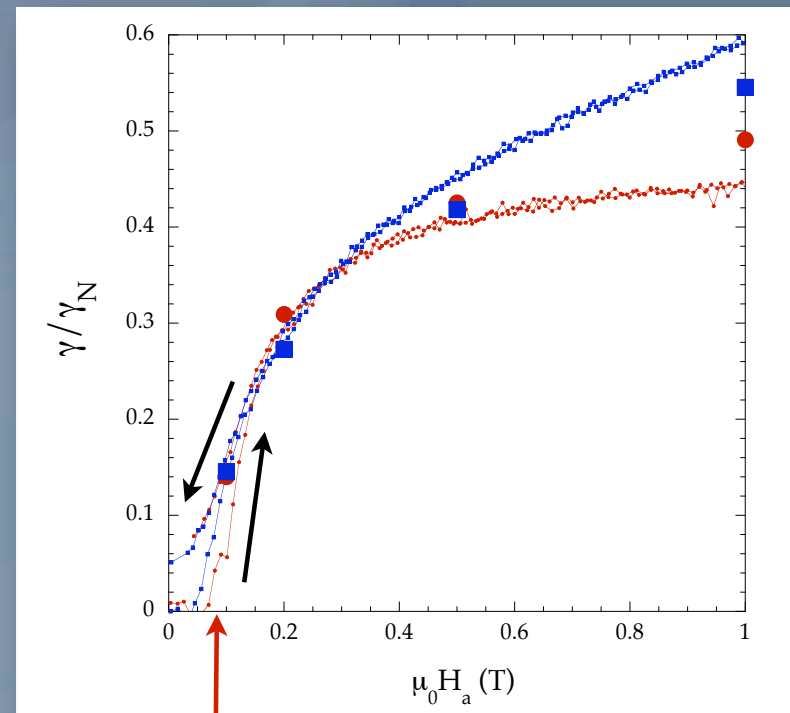
Linear field dependence for $H < 0.2T$
but vicinity of $H_p \sim 500-1000G$?
isotropic ($H < 0.5-1T$)

starting point : detailed study of the
field dependence at low B
is γ linear in B at low B or
sublinear down to $B = 0$?

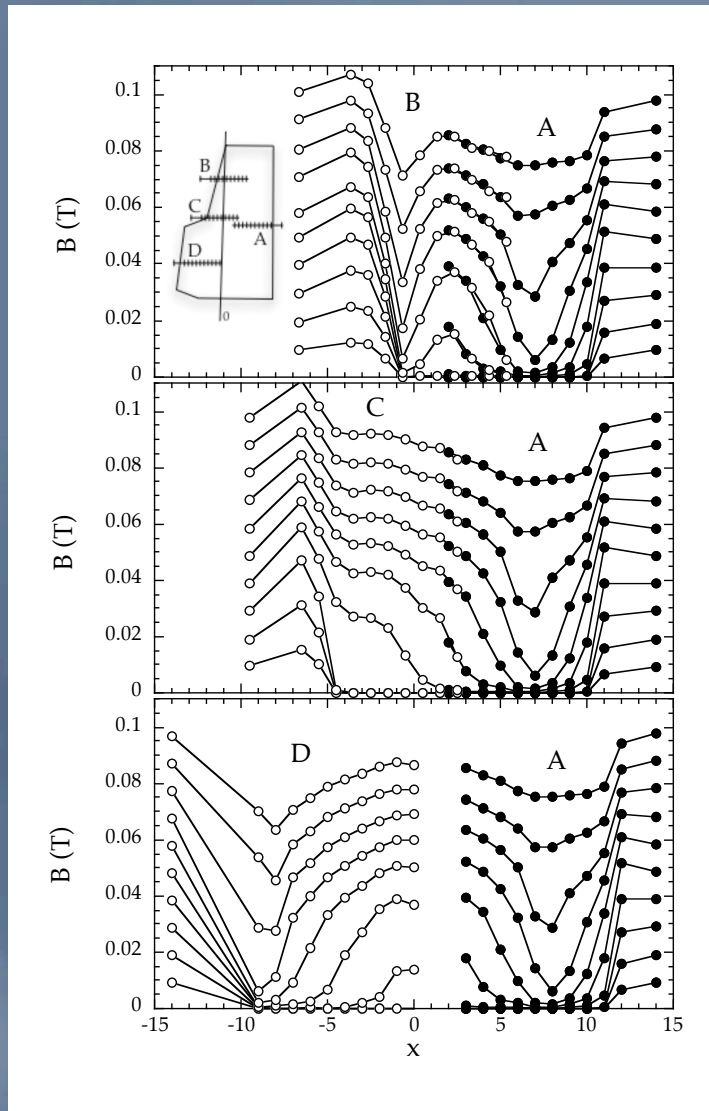
AC technique



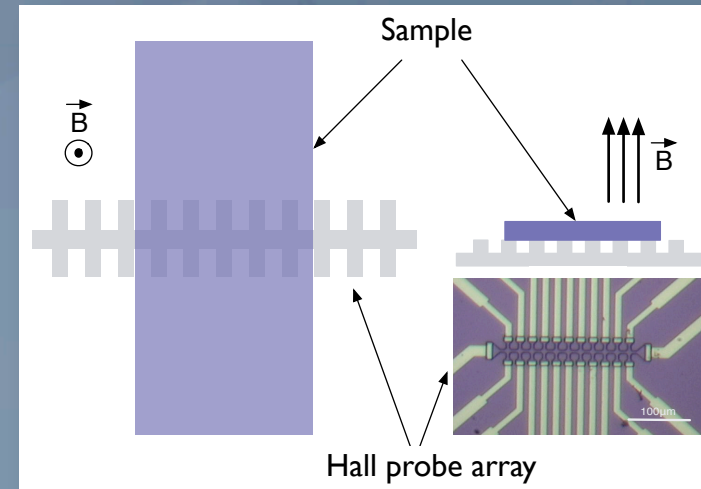
high sensitivity
well adapted for small samples
continuous field measurements



H_p

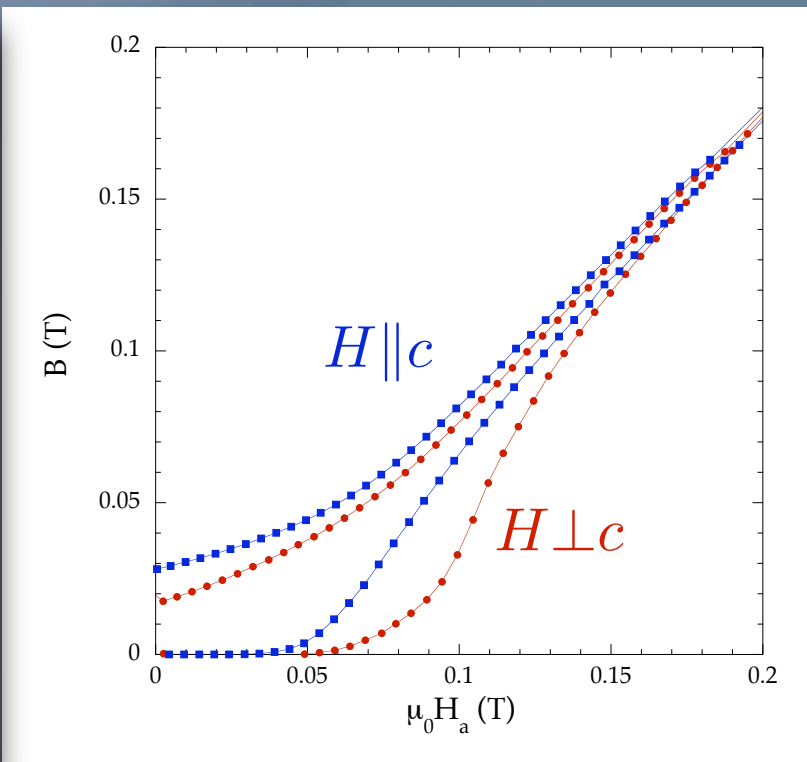
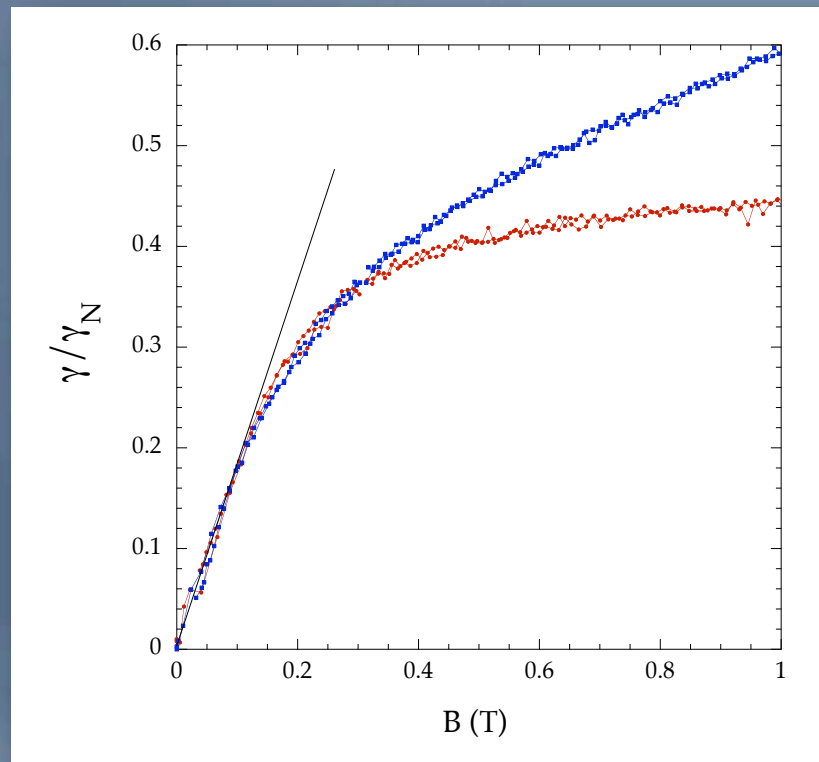


Hall probe magnetometry



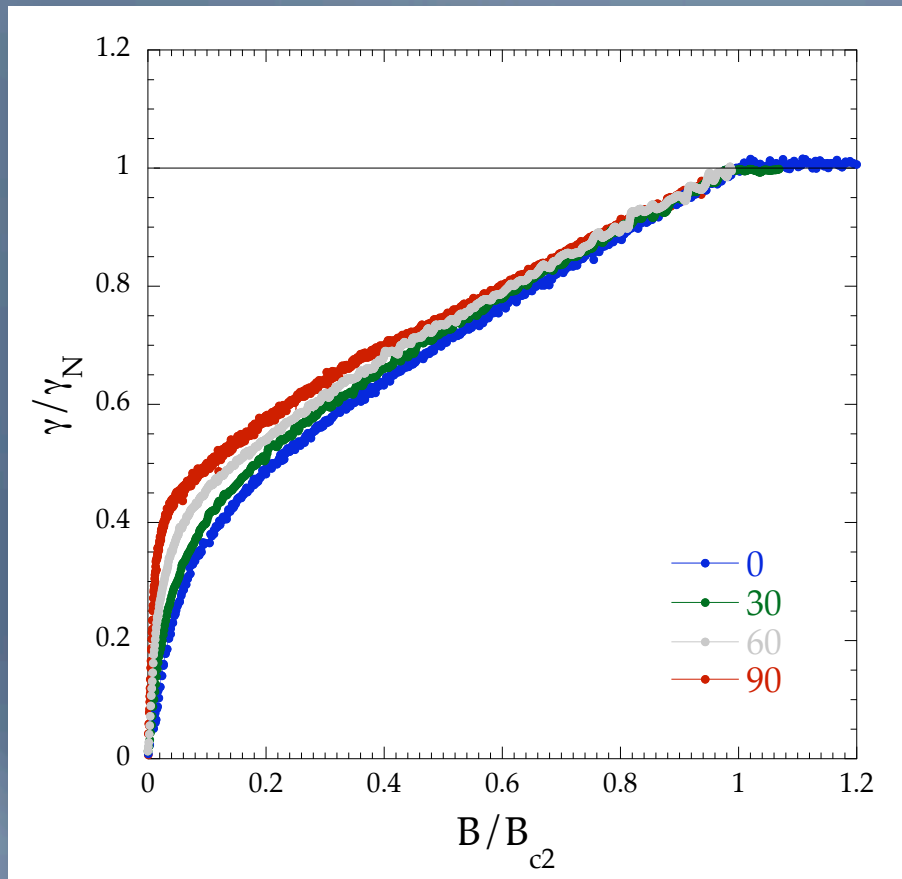
$$B^{(av)} = \frac{\sum B_n}{n}$$

Dome shape B-profiles :
influence of geometrical barriers



γ linear only for $B < 800\text{G}$

isotropic up to $\sim 0.3\text{T}$



high field

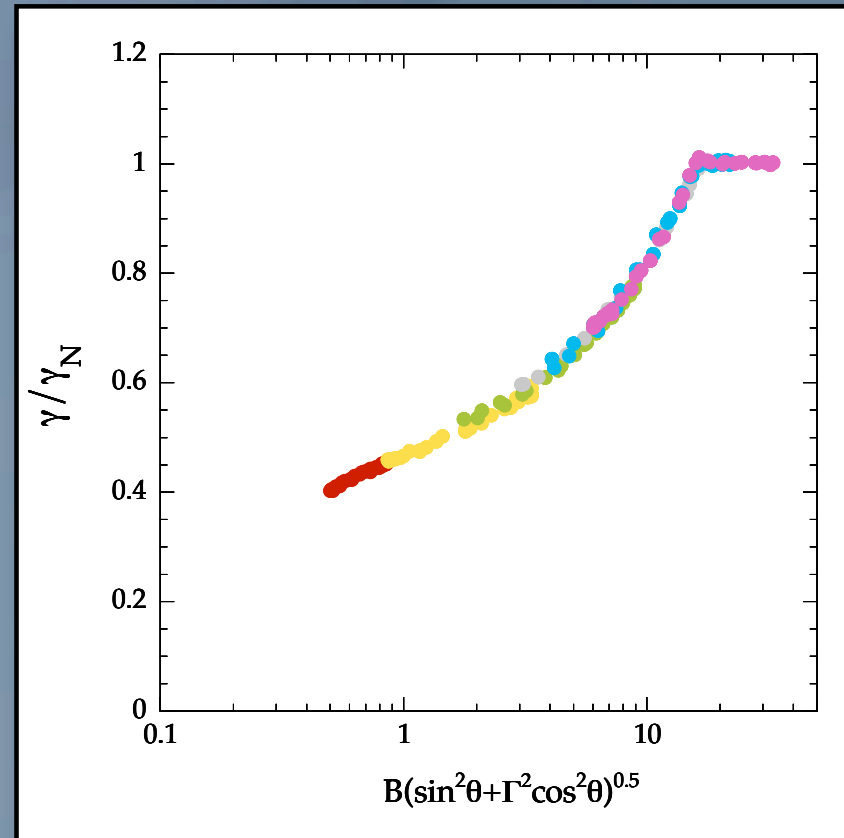
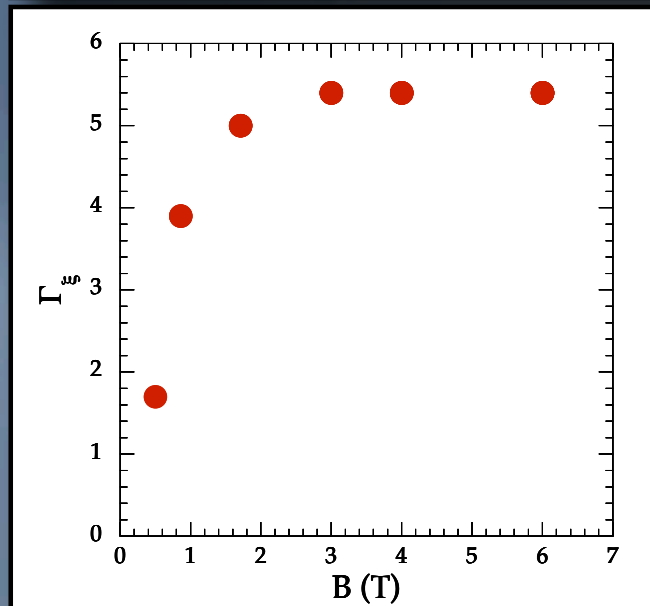
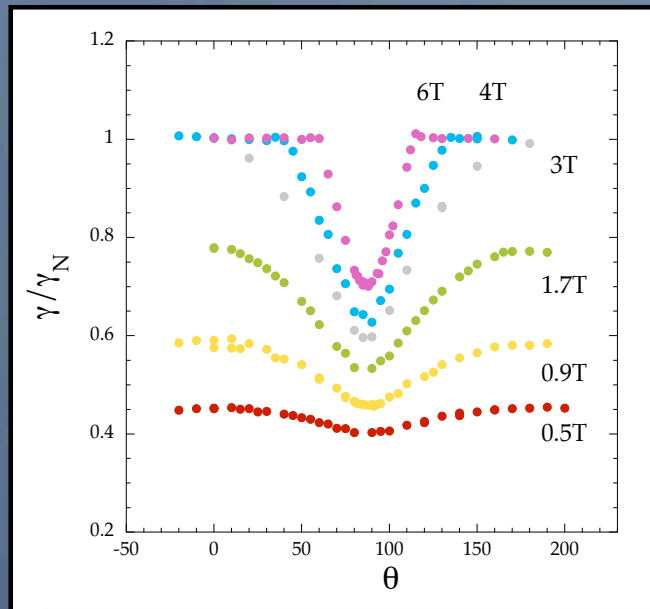
$$\gamma = f\left(\frac{B}{B_{c2}}\right)$$

$$B_{c2}(\theta) = \frac{B_{c2}^{ab}}{\sqrt{\sin^2\theta + \Gamma_\xi^2 \cos^2\theta}}$$

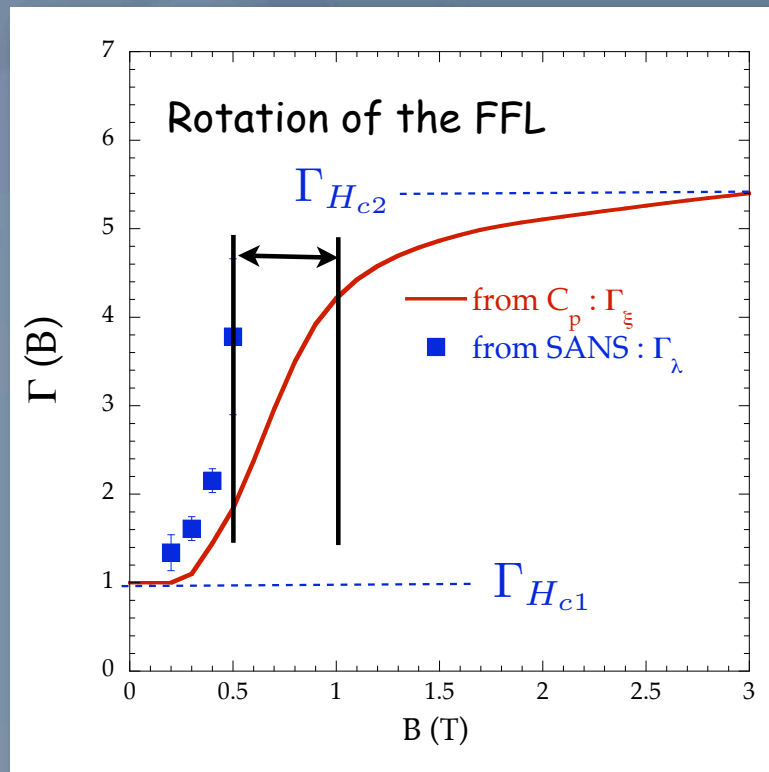
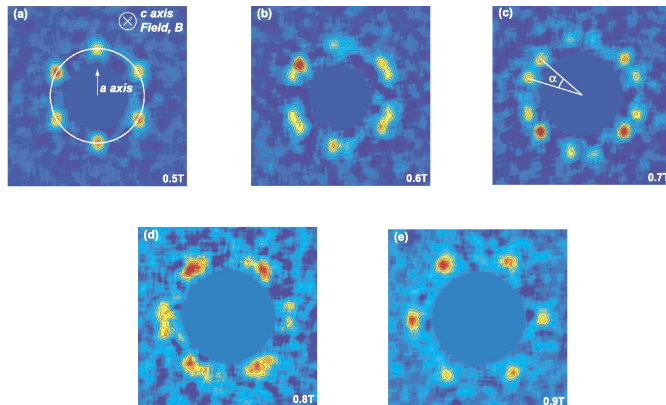
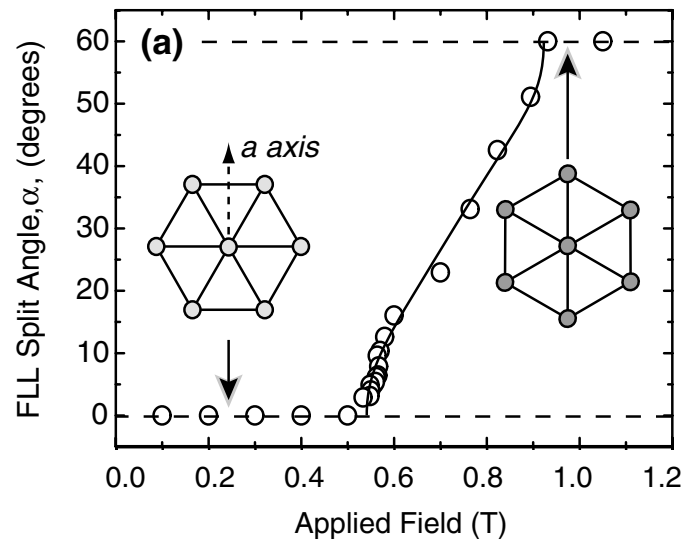
$$\Gamma \sim 5.4$$

but at low field γ is isotropic ($\Gamma \sim 1$).....

$$B_{c2}(\theta) = \frac{B_{c2}^{ab}}{\sqrt{\sin^2\theta + \Gamma_{\xi}^2(B)\cos^2\theta}}$$



$$\gamma_{H\parallel c}(\Gamma_{\xi}(B)B) = \gamma_{H\perp c}(B) = \gamma_{\theta}(B\sqrt{\sin^2\theta + \Gamma_{\xi}^2(B)\cos^2\theta})$$

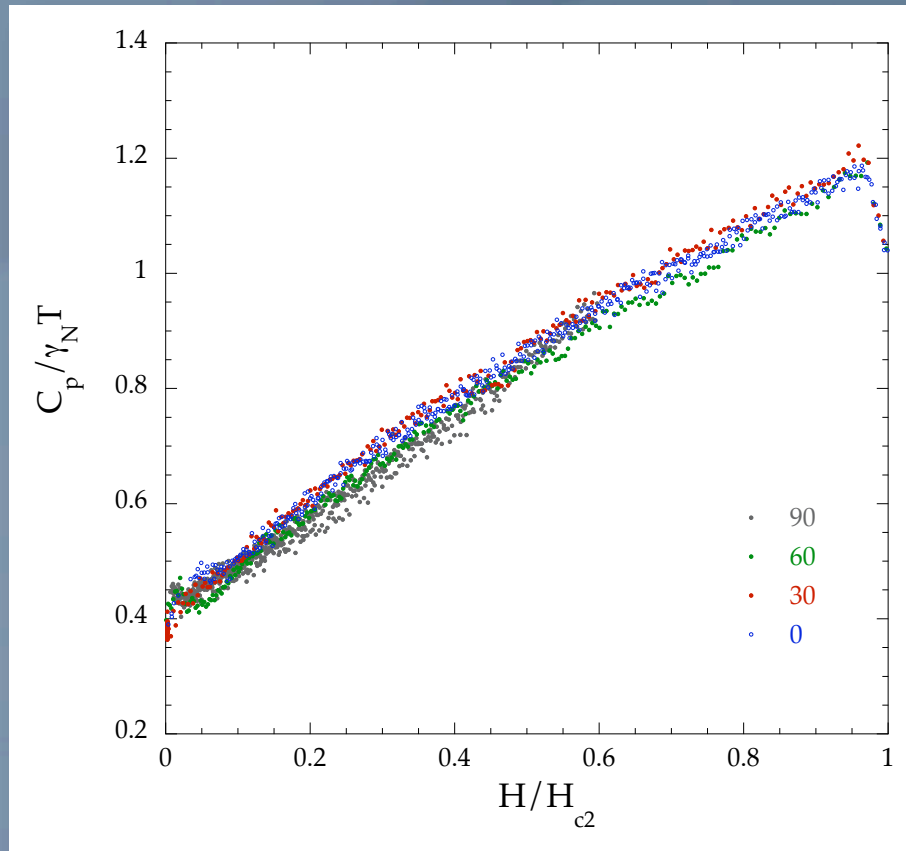


$$\Gamma_{H_{c1}} \leq \Gamma_{\lambda}(B) = \Gamma(B) = \Gamma_{\xi}(B) \leq \Gamma_{H_{c2}}$$

$$T=9\text{K} > \Delta_{\pi}$$

the low field
isotropic behaviour
is not present anymore

$$C_p = f(H/H_{c2})$$



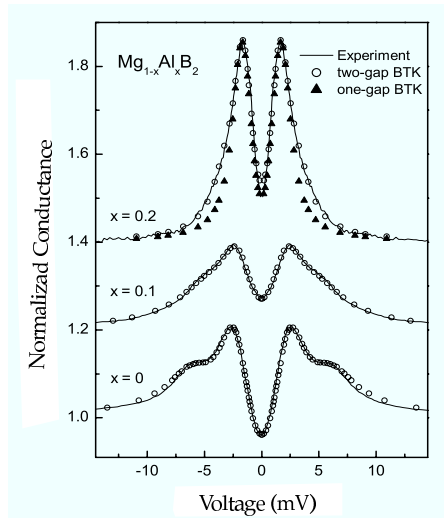
The background is a solid blue color with several large, overlapping, wavy, and translucent shapes that create a sense of movement and depth. These shapes are lighter blue and white, blending into the darker blue background.

Aluminum doping....

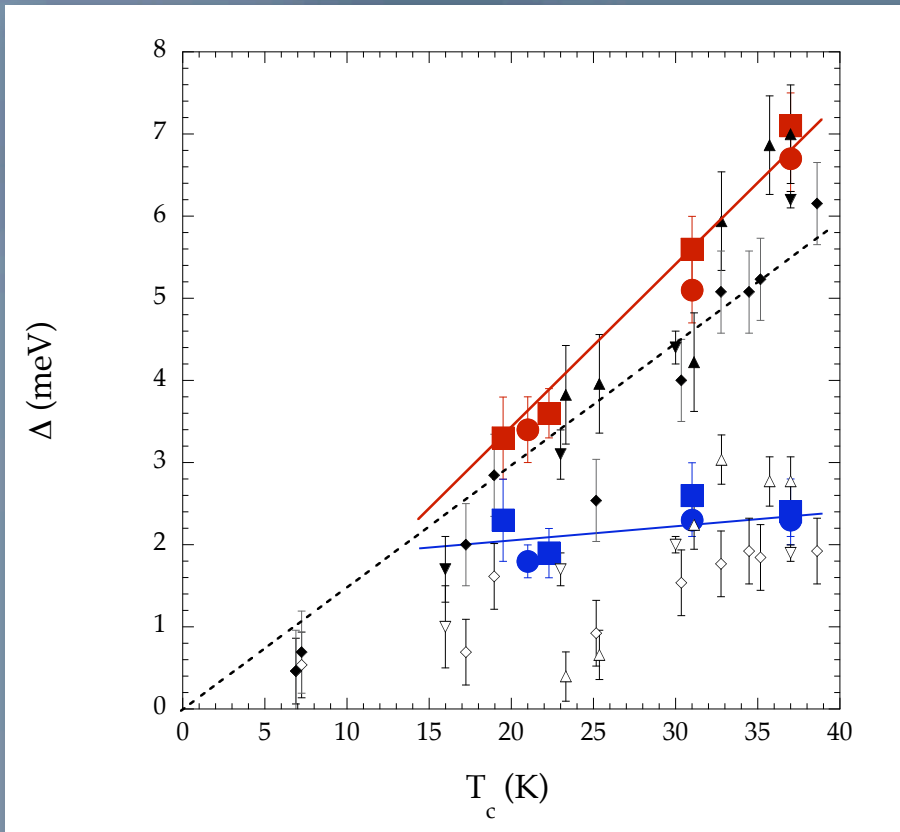


Electronic doping

$$C_p = \omega_\sigma C_p^\sigma(T, \Delta_\sigma/T_c) + (1 - \omega_\sigma) C_p^\pi(T, \Delta_\pi/T_c)$$



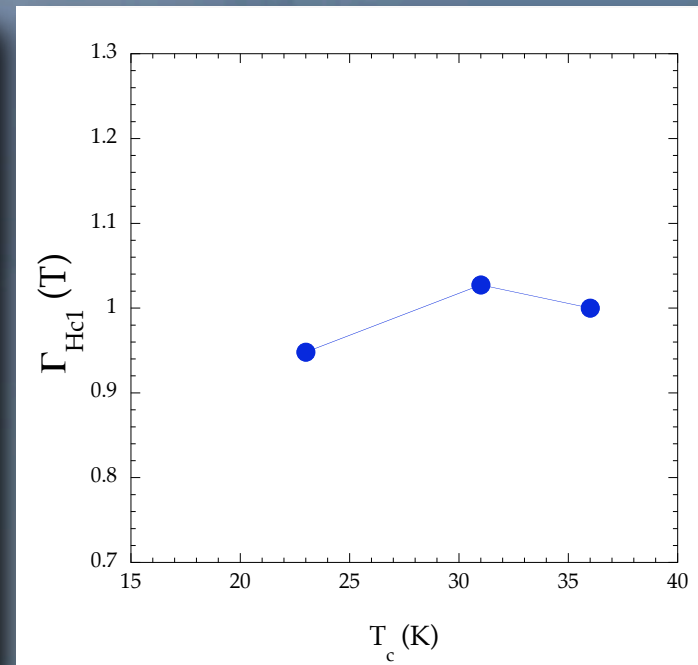
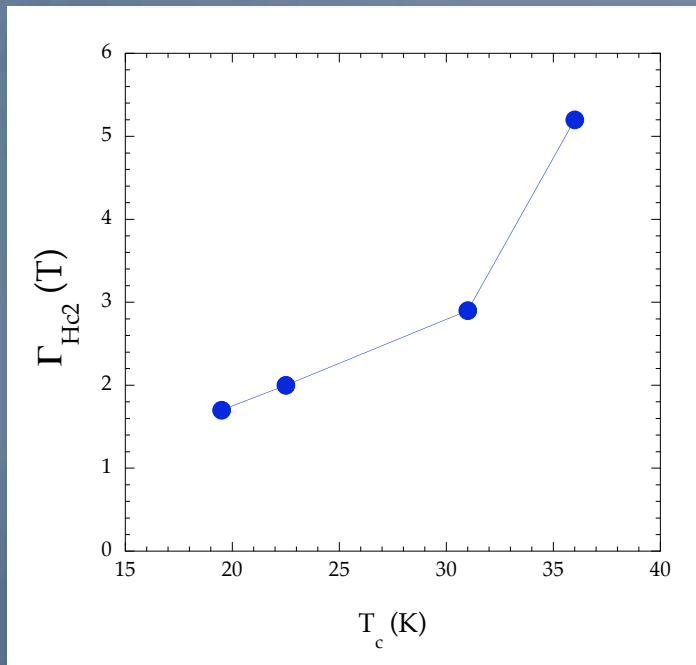
coll. P.Samuely et al. SAS, Kosice



small gap \sim constant

increase of the
intraband scattering

possible merging
of the gaps for $T_c \sim 10\text{K}$



$$H_{c2}^c \propto \frac{\Delta_\sigma^2}{v_{F,ab}^2} \quad \Delta_\sigma \downarrow, v_{F,ab} \downarrow, H_{c2}^c \sim \text{cte}$$

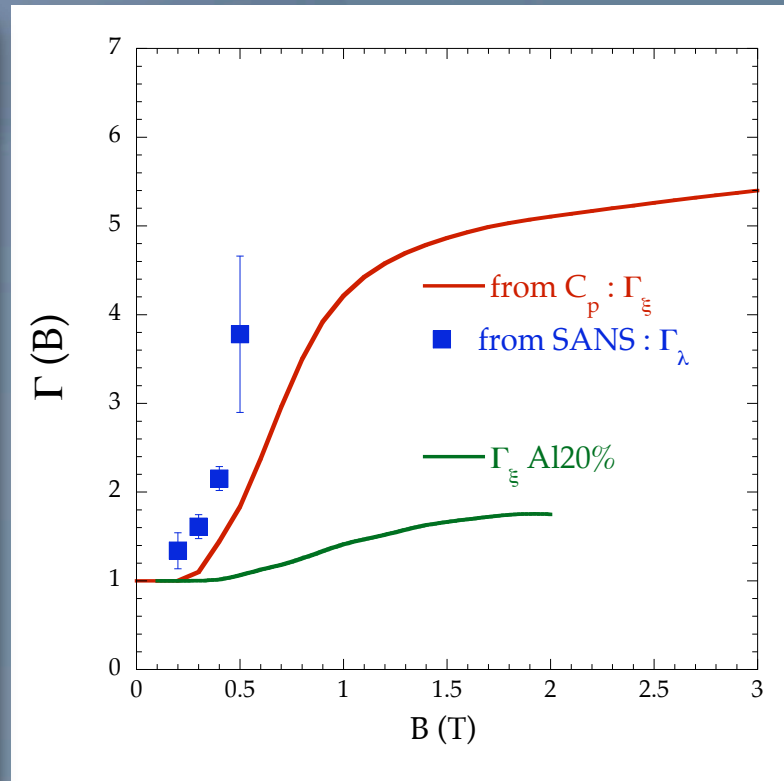
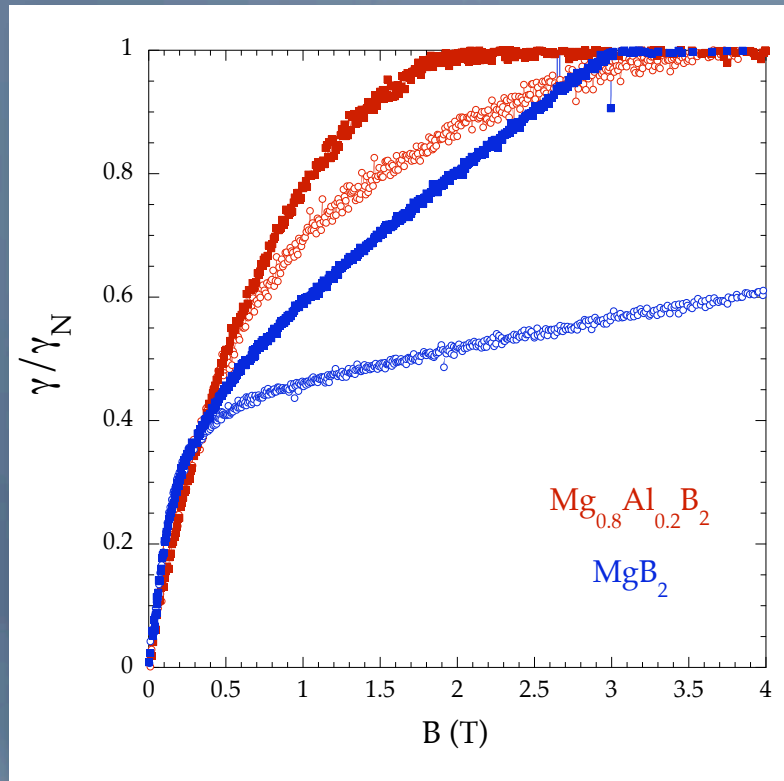
$$H_{c2}^{ab} \propto \frac{\Delta_\sigma^2}{v_{F,ab} v_{F,c}} \quad v_{F,c} \sim \text{cte}, H_{c2}^{ab} \downarrow$$

superconducting density \downarrow

$$H_{c1} \downarrow$$

$$\Gamma_{H_{c2}} = \frac{v_{F,ab}}{v_{F,c}} \quad \downarrow$$

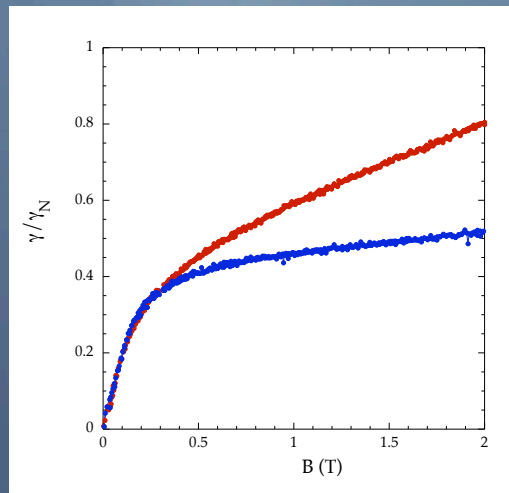
$$\Gamma_{H_{c1}} \approx 1$$



towards an isotropic one gap superconductor

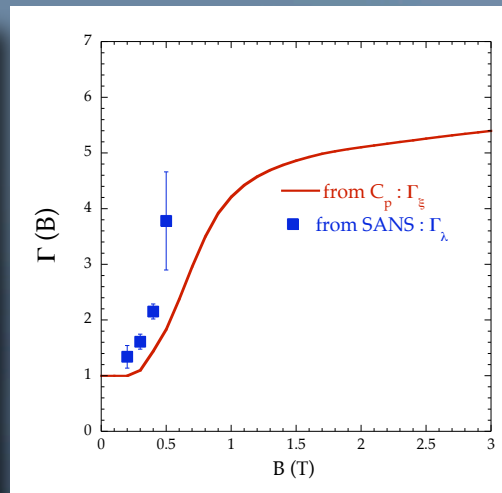
Conclusions

Linear B dependence
for $B < 800\text{G}$
getting strongly sublinear
for higher fields



Field dependent
anisotropy

$$C_p = f(H/H_{c2}) \text{ for } T > \Delta_\pi$$



Al doping :
towards an isotropic
one gap superconductor

