

# Superconductivity in $\text{MgB}_2$ and B-doped diamond

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- two gap superconductivity in  $\text{MgB}_2$

L. Lyard, C. Marcenat, J. Marcus

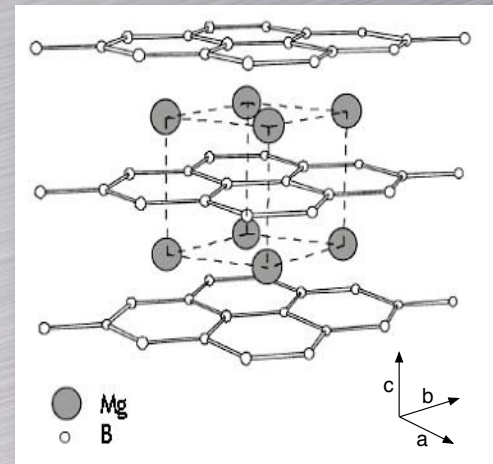
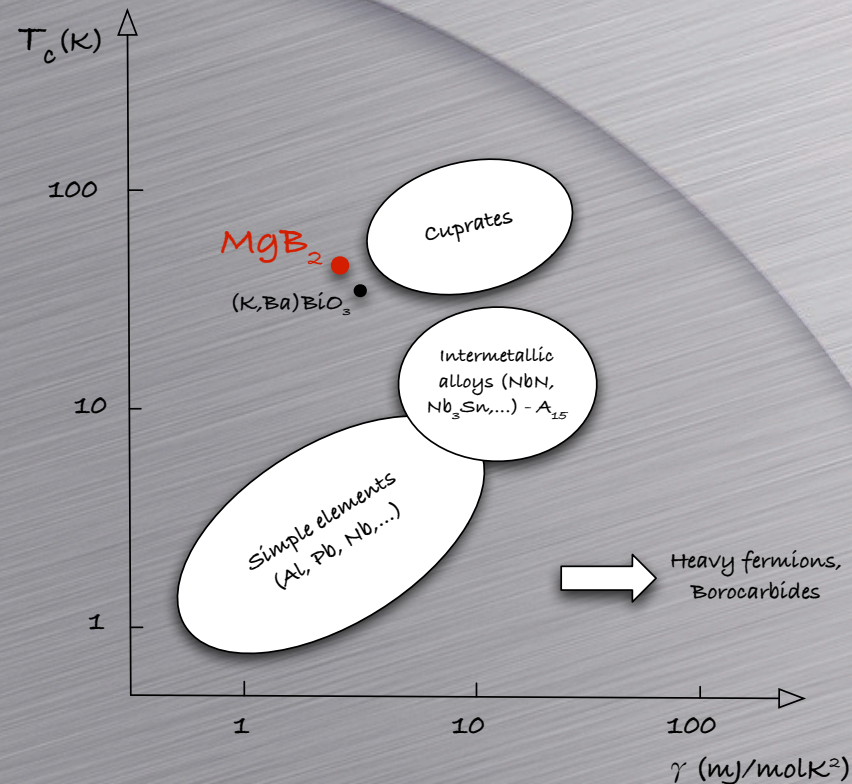
- superconductivity in B-doped diamond : proximity of a metal-insulator transition

E. Bustarret, E. Gheeraert, C. Marcenat, F. Gustafsson, J. Marcus



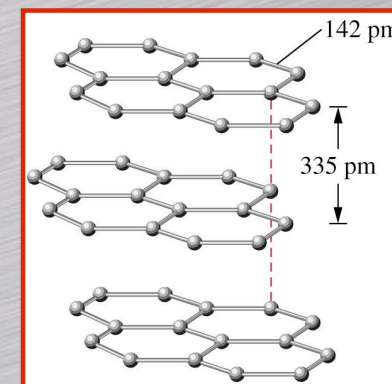
# MgB<sub>2</sub> : a new "high T<sub>c</sub>" superconductor

T<sub>c</sub> ~ 40K



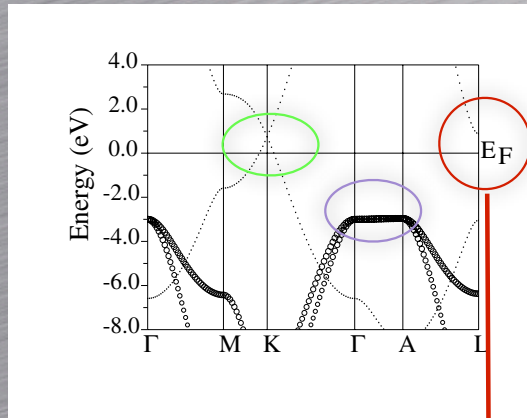
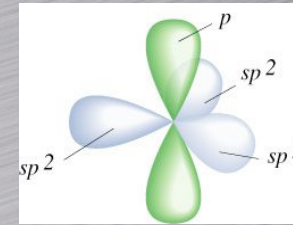
hexagonal structure

B planes similar to C planes in graphite



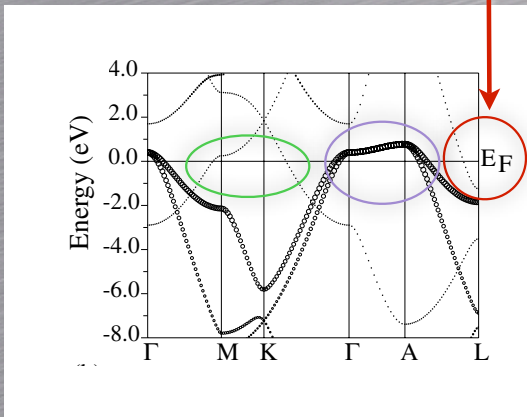


- 2 kinds of electronic orbitals derived from
- in plane  $sp_2$ -boron orbitals :  $\sigma$  bands
  - $p_z$  boron orbitals :  $\pi$  bands



"same" electronic structure than in Graphite  
in which the  $\sigma$  bands are full  
(involved in covalent bounds)

→  $\pi$  band : small e/hole pockets close to the K point



in  $MgB_2$  the "second"  $\pi$  band is lowered  
by attractive  $Mg^{2+}$  potential

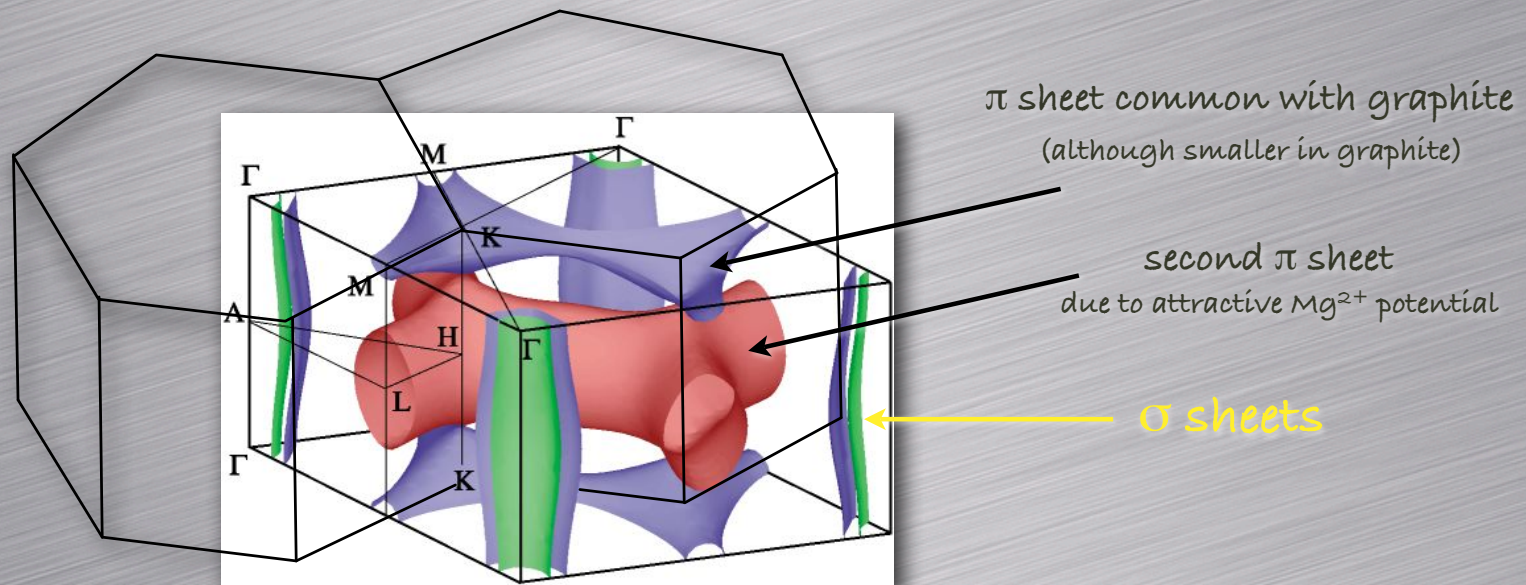
charge transfert from  $\sigma$  to  $\pi$  bands

→ unfilled  $\sigma$  bands



"2" Fermi surface sheets

actually  $2 \times \pi$  (3D) and  $2 \times \sigma$  (quasi-2D) bands



the quasi 2D characters leads to a "large" DOS  $\sim 0.30$  st/eV.cell

(close to the 2D value  $m^* a^2 / \pi \hbar^2 \sim 0.33$  st/eV.cell)

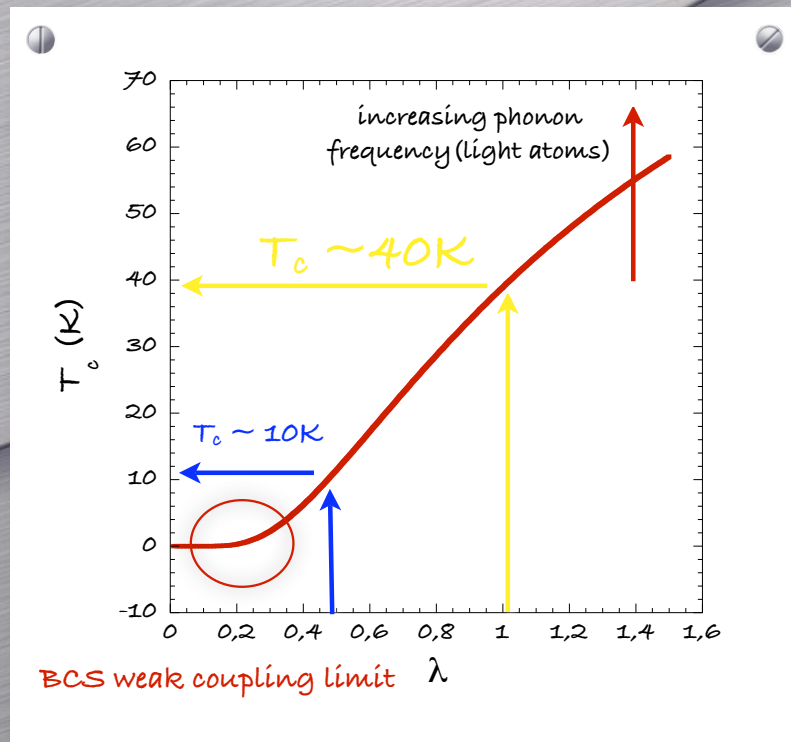
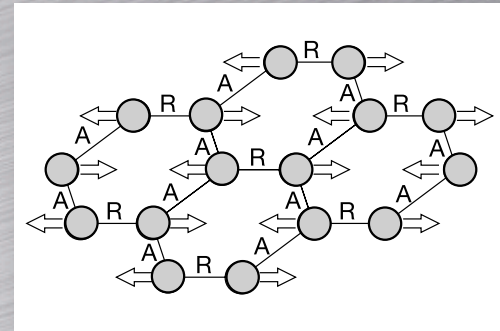
despite the small hole doping level  $\sim 0.07$  hole/B



strong coupling of  $\sigma$  electrons with  $E_{2g}$  vibration mode  
+ large DOS

-> electron-phonon coupling constant  $\lambda \sim 1.0$

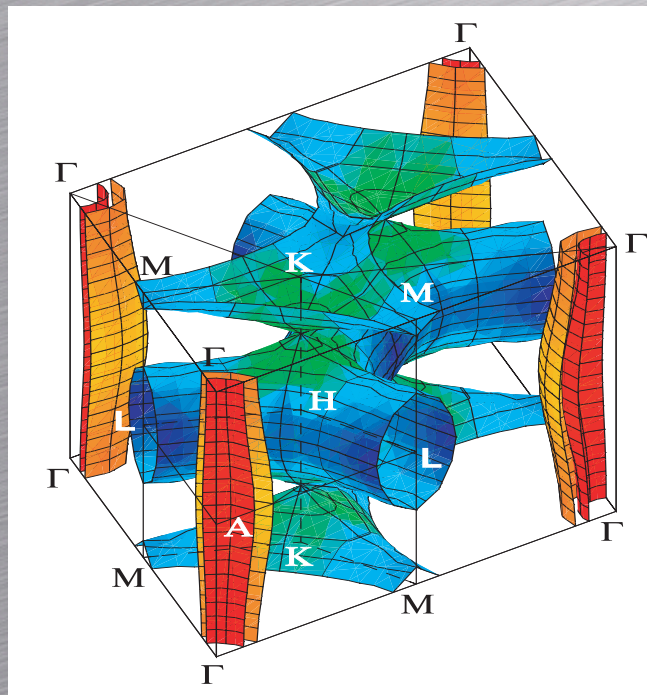
despite a larger DOS ( $\sim 0.4 \text{ st/eV.cell}$ )  
the coupling is much worse in the  $\pi$  band ;  $\lambda \sim 0.45$



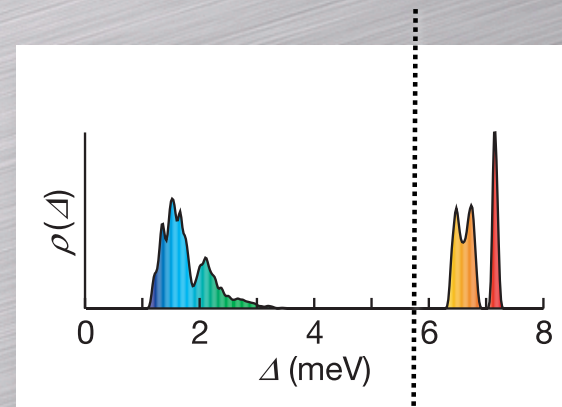
Mac-Millan expansion  
screening coefficient : 0.15  
 $\omega = 540K$



Two co-existing superconductors ????



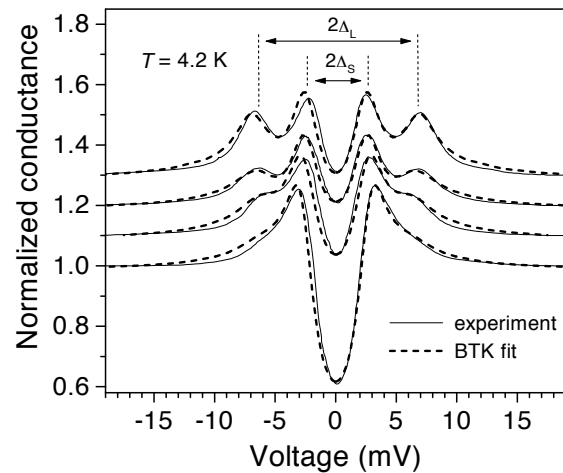
2 SF sheets  $\longrightarrow$  2 gaps



BCS value

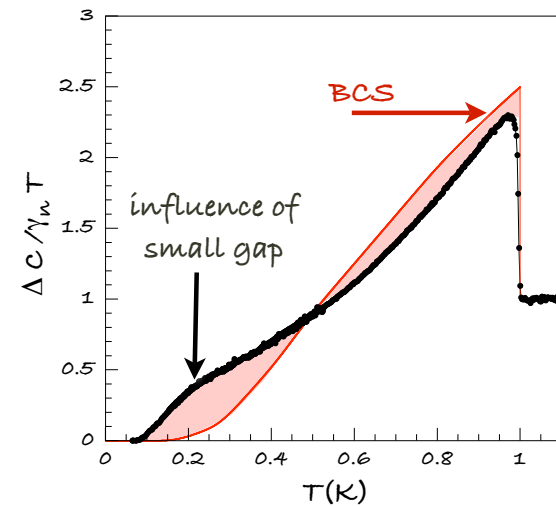


## point contact spectroscopy



P. Szabo et al. PRL 01

confirmed by specific heat  
(i.e. bulk) measurements

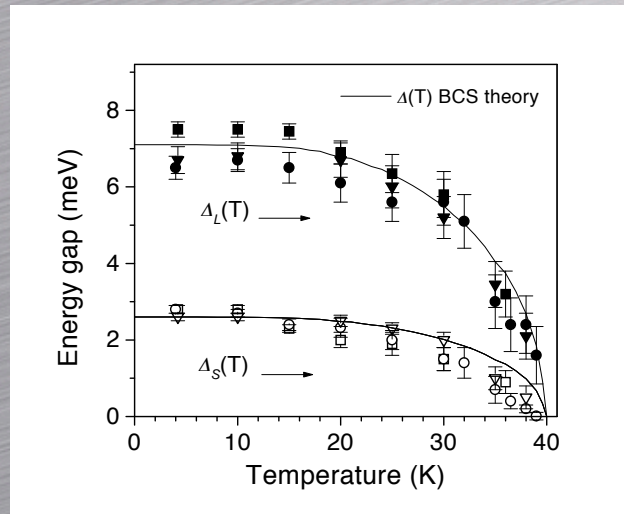


see also F. Bouquet et al. EPL 01

$$C_p/T : \gamma \propto e^{-\Delta/kT}$$

2 gaps  $\Rightarrow$  2  $T_c$  values ???

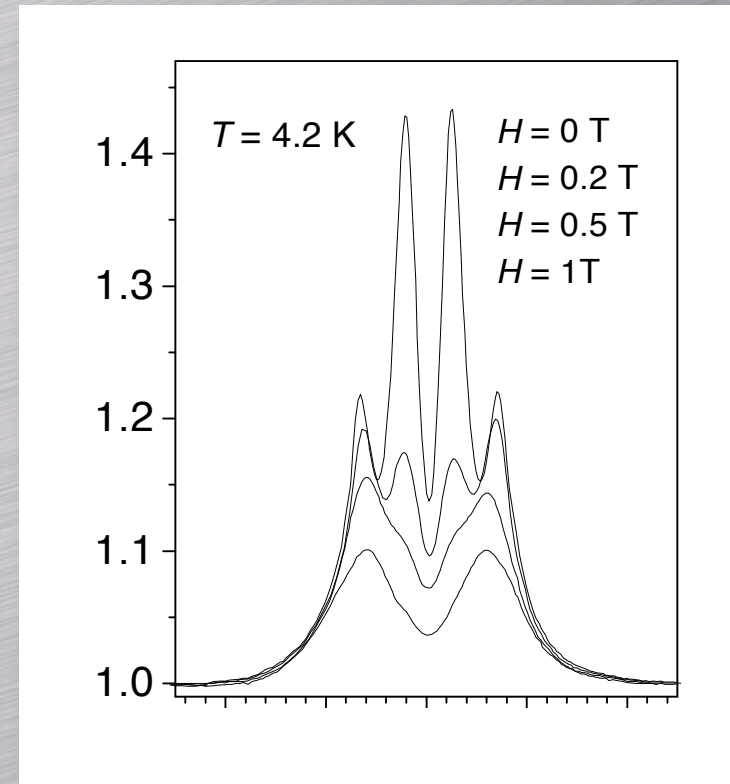
but present different field dependences



both gaps are closing at the **same**  $T_c$   
 $\Rightarrow$  weakly coupled bands

zero coupling  
 $\Rightarrow 2 T_c$ 's

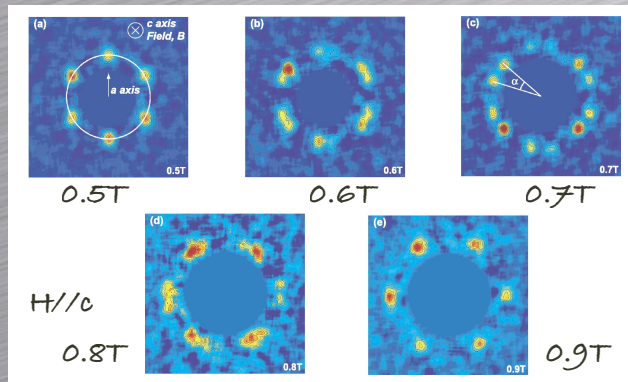
strong coupling  
 $\Rightarrow$  FS averaged  
 $\lambda$  value ( $\sim 0.7$ )  
 $\Rightarrow T_c \sim 20K$



the small gap is very  
sensitive to  $H$  i.e.  
"disappears" around 1T

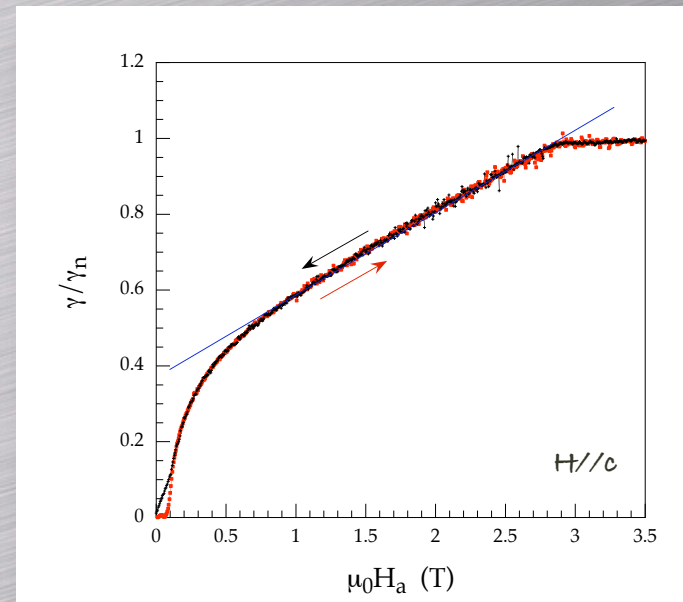


and the vortex lattice  
rotates by  $30^\circ$



SANS measurements R.Cubitt et al. PRL 03

+ anomalous field dependence of  $\gamma$



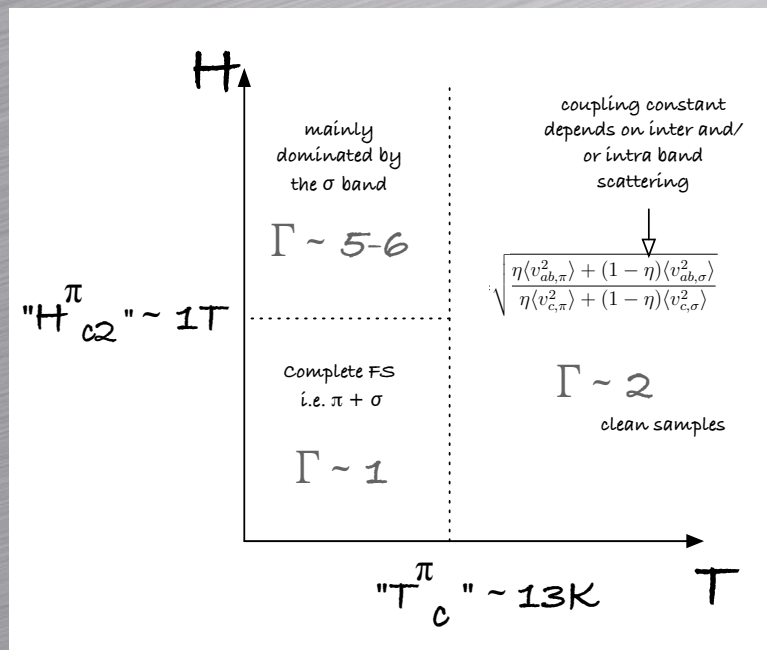
see also F.Bouquet et al. PRL 02

Should the system be described by two  $\xi$  and two  $\lambda$  values (associated with each band) or only one field dependent value (as we believe)?  
In any case the "properties" of the superconducting state are going to be strongly field dependent: *exemple of the anisotropy*



$$\Gamma = \frac{\langle m_{ab}^* \rangle_{FS}}{\langle m_c^* \rangle_{FS}} = \sqrt{\frac{\langle v_{F,c}^2 \rangle_{FS}}{\langle v_{F,ab}^2 \rangle_{FS}}} = \frac{\lambda_c}{\lambda_{ab}} = \frac{\xi_{ab}}{\xi_c}$$

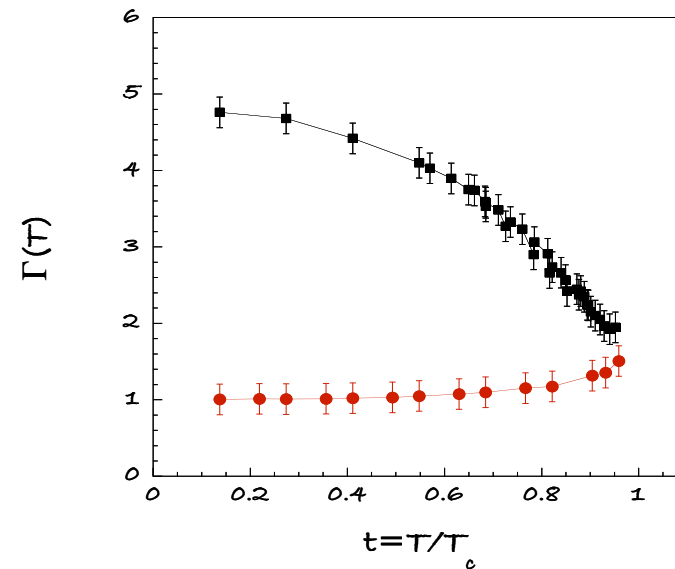
Which Fermi Surface should we take into account ???



How can we get  $\Gamma$  ???

high field :  $\Gamma = \Gamma_{H_{c2}} = \frac{H_{c2,ab}}{H_{c2,c}}$

low field :  $\Gamma \sim \Gamma_{H_{c1}} = \frac{H_{c1,c}}{H_{c1,ab}}$



from  $C_p$  and magnetotransport, Lyard et al. PRB 02

from  $C_p$  and Hall probe magnetometry, Lyard et al. PRL 04



what about Graphite.....

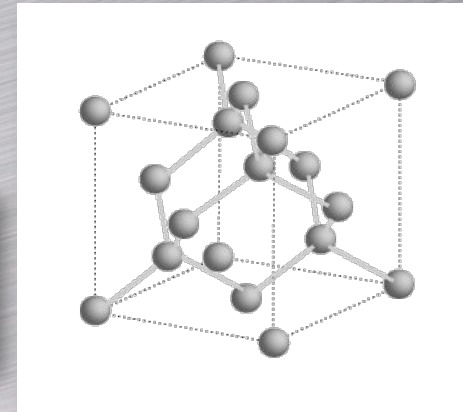
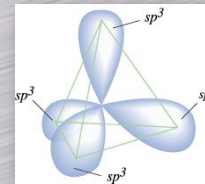
doping by intercalation (Na,...) leads to  
superconductivity :  $T_c \sim 5K$

but superconductivity  
remains in the "soft"  $\pi$  band  
(similar to alkali doped fullerenes)



DIAMOND  
very "strong"  $\sigma$  bounds

but 3D ( $sp^3$ )....



Superconductivity ( $T_c \sim 7K$ ) has been  
observed in Ba doped Si-clathrates  
(cage like structure with  $sp^3$  bounds)



all 4e are involved in covalent bounds -> large gap semiconductor

can be either n or p doped by substituting C atoms by P or B atoms, respectively

In B-doped samples, the system becomes **metallic** when the boron impurity band overlaps the diamond valence band  
i.e. for boron concentrations > a few  $10^{20} \text{ cm}^{-3}$

what about the e-phonon coupling  $\lambda$  ???

$$\lambda = N(E_F) \cdot \frac{I^2}{M\omega^2}$$

DOS  
st/ev.spin."2 atoms-cell"

FS averaged e-ion  
matrix element

"appropriately defined"  
mean square frequency

$\pm u$  : splitting of the top of the valence band  
for the "appropriate" displacement  $u$

	N	I (eV/Å)	$\omega(\text{cm}^{-1})$	$I^2/M\omega^2$	$\lambda$
MgB <sub>2</sub>	0.15	12	540	6.7	1.0
C-B3%	$0.06 \pm 0.01$	$24 \pm 3$	1080	$7.5 \pm 0.8$	$0.45 \pm 0.1$

2x the already large  
MgB<sub>2</sub> value

1/2 of MgB<sub>2</sub> value to  
smaller DOS (3D)



	method	model	B/C at.%	$n_B$ $10^{20} \text{ cm}^{-3}$	e/ph coupling $\lambda$	$T_c$ (K)	Remark
Boeri et al PRL 93, 237002	First Principles LMTD	VCA Virtual crystal	3	50	0.37	0.2	$T_c = 25 \text{ K}$ for 10 at%
Lee et al PRL 93, 237003	First principles APW-CPA	VCA Virtual crystal	2.5	44	0.53	9	$E_F$ at $\text{VBM}^u - 0.6 \text{ eV}$
Xiang et al Cond-mat 0406446	First principles supercell DFT-LDA	$C_{35}B$	2.8	49	0.39	4.4	B modes contribute to e/ph coupling
Blase et al PRL 93, 237004	Ab initio Supercell DFT-LDA	$C_{53}B$	1.85	33	0.43	4	« Local » C-B modes couple with free carriers $E_F$ at -0.52 eV

phonon modes involved in e/ph  
coupling not clearly identified

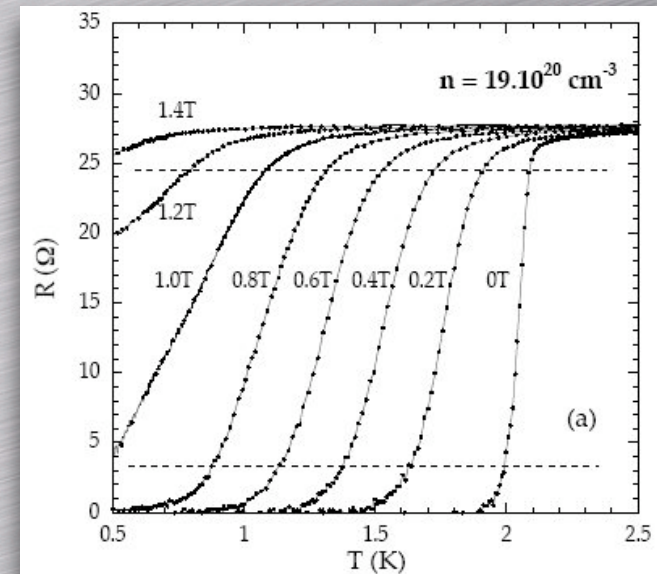
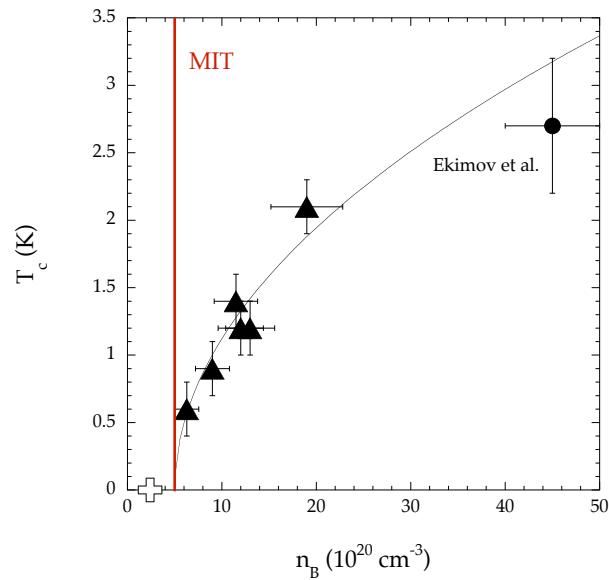


Ekimov et al. Nature 04

- Polycrystalline diamond (HPHT bulk)
- Doping level :  $10^{21}$  B/cm<sup>3</sup> :  $T_c \sim 3K$

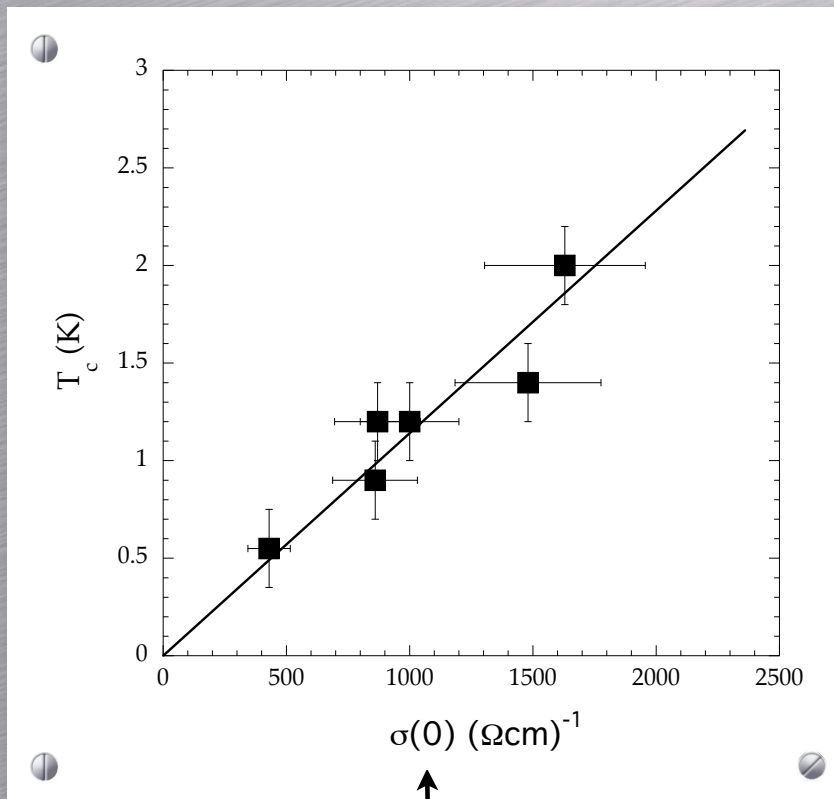
confirmed on homoepitaxial films

E. Bustarret 04





is there any relation between this metal-insulator transition and superconductivity ???



$$T_c \sim \sigma(0) ???$$

enhanced superconductivity  
due to reduced screening close  
to the MIT ???

Soulen - Osofsky et al.

normal state conductivity extrapolated to zero



Hole doping of the  $\sigma$  bands leads to very efficient electron-phonon interaction potential



+ large DOS due to 2D character ( $sp^2$ ) in  $MgB_2$   
 $\Rightarrow T_c \sim 40K$



superconducting is also "induced" in the  $\pi$  band  
 $\rightarrow$  two gap superconductivity



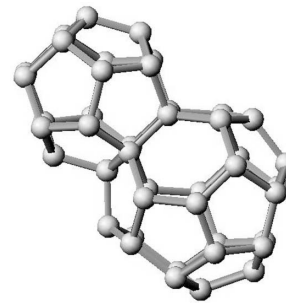
anomalous field and temperature dependence of the superconducting properties



can explain the onset of superconductivity in B-doped diamond in, which "reduced" DOS  $\Rightarrow T_c \sim$  a few K



superconductivity appears in the vicinity of a MIT



large  $T_c$  values predicted in C clathrates ( $X_8@C_{46}$ )

$\lambda \sim 1.4$  &  $\omega \sim 1500K$

$\Rightarrow T_c$  between 50

and ...150K !!!

(depending on screening coefficient)  
Connétable et al. 2003